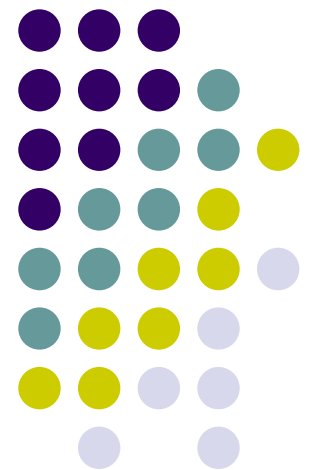
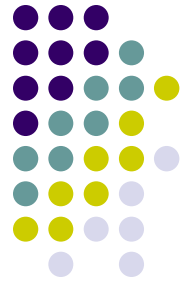


5. Transformasi Integral dan Persamaan Integral

- 5.1. Transformasi Integral
- 5.2. Transformasi Laplace
- 5.3. Transformasi Fourier
- 5.4. Persamaan Integral





5.1. Transformasi Integral

Di dalam Fisika Matematika kita sering menjumpai pasangan fungsi yang dihubungkan sbb:

$$g(\alpha) = \int_a^b f(t)K(\alpha, t)dt$$

Fungsi $g(\alpha)$ disebut transformasi (integral) dari $f(t)$ dengan kernel $K(\alpha, t)$.

→ Mapping fungsi $f(t)$ di ruang- t dari fungsi $g(\alpha)$ di ruang- α .



Mengapa kita butuh transformasi integral??

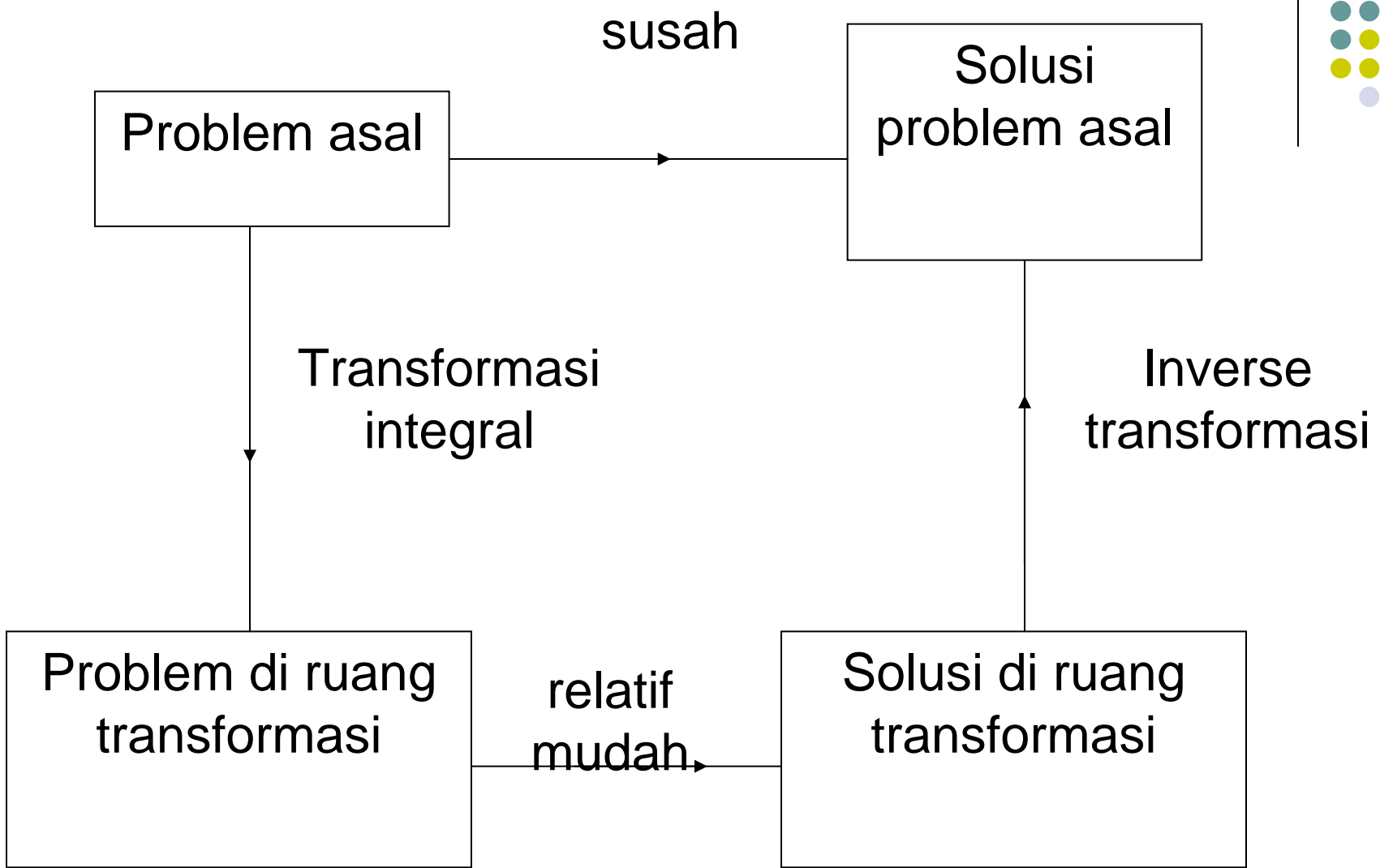
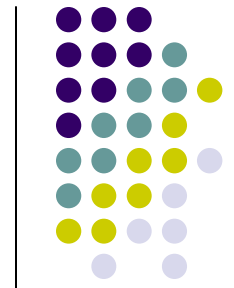
Karena dalam banyak kasus masalah lebih mudah diselesaikan dengan cara transformasi dan inversi.

Kita membutuhkan representasi dalam ruang lain.

Contoh di fisika:

waktu \Leftrightarrow frekuensi

ruang real \Leftrightarrow ruang momentum



Satu diantara transformasi yang terpenting → Fourier

$$g(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t)e^{i\alpha t} dt$$

Ada tiga lainnya:

Transformasi Laplace:

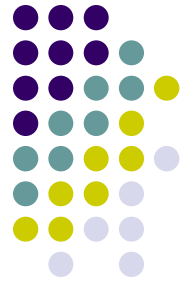
$$g(\alpha) = \int_0^{\infty} f(t)e^{-\alpha t} dt$$

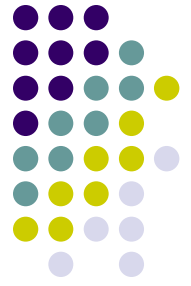
Transformasi Hankel (Fourier-Bessel)

$$g(\alpha) = \int_a^b f(t)tJ_n(\alpha t)dt$$

Transformasi Mellin

$$g(\alpha) = \int_a^b f(t)t^{\alpha-1} dt$$





5.2. Transformasi Laplace

Transformasi Laplace $f(s)$ atau L dari $F(t)$ didefinisikan:

$$f(s) = L\{F(t)\} = \lim_{a \rightarrow \infty} \int_0^a e^{-st} F(t) dt = \int_0^{\infty} e^{-st} F(t) dt$$

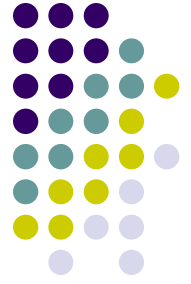
Beberapa fungsi sederhana:

1) $F(t) = 1, t > 0$

$$L\{1\} = \int_0^{\infty} e^{-st} dt = \frac{1}{s}$$

2) $F(t) = e^{kt}, t > 0$

$$L\{e^{kt}\} = \int_0^{\infty} e^{-st} e^{kt} dt = \frac{1}{s-k}, \text{ untuk } s > k$$



3) Fungsi hiperbolik sinus dan kosinus

Karena:

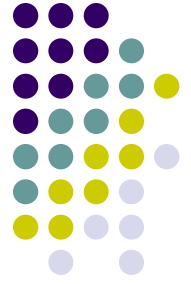
$$\cosh(kt) = \frac{1}{2}(e^{kt} + e^{-kt}) \quad \text{dan}$$
$$\sinh(kt) = \frac{1}{2}(e^{kt} - e^{-kt})$$

Maka

$$L\{\cosh(kt)\} = \frac{1}{2} \left(\frac{1}{s-k} + \frac{1}{s+k} \right) = \frac{s}{s^2 - k^2}$$

dan

$$L\{\sinh(kt)\} = \frac{1}{2} \left(\frac{1}{s-k} - \frac{1}{s+k} \right) = \frac{k}{s^2 - k^2}$$



4) Fungsi sinus dan kosinus biasa

dengan menggunakan:

$$\cos(kt) = \cosh(ikt) \text{ dan } \sin(kt) = -i \sinh(ikt)$$

diperoleh:

$$L\{\cos(kt)\} = \frac{s}{s^2 + k^2}$$

dan

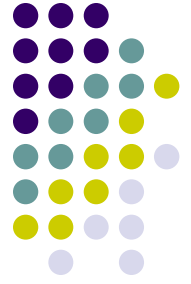
$$L\{\sin(kt)\} = \frac{k}{s^2 + k^2}$$

5) $F(t) = t^n$

$$L\{t^n\} = \int_0^{\infty} e^{-st} t^n dt = \frac{n!}{s^{n+1}}$$

Tabulasi:

$F(t)$	$f(s)$
1	$\frac{1}{s}$
e^{kt}	$\frac{1}{s - k}$
$\cosh(kt)$	$\frac{s}{s^2 - k^2}$
$\sinh(kt)$	$\frac{k}{s^2 - k^2}$
$\cos(kt)$	$\frac{s}{s^2 + k^2}$
$\sin(kt)$	$\frac{k}{s^2 + k^2}$
t^n	$\frac{n!}{s^{n+1}}$



Tabel lengkap
dapat dilihat di
Arfken



Contoh soal:

1. Carilah $F(t)$ bila

$$f(s) = \frac{k^2}{s(s^2 + k^2)}$$

Jawab:

Fungsi $f(s)$ dapat diuraikan menjadi:

$$f(s) = \frac{A}{s} + \frac{Bs + C}{(s^2 + k^2)}$$

dengan substitusi balik, diperoleh $A = 1$, $B = -1$, $C = 0$, sehingga:

$$f(s) = \frac{1}{s} - \frac{s}{(s^2 + k^2)}$$

dengan demikian inverse $f(s)$ menjadi:

$$F(t) = 1 - \cos(kt)$$



2. Carilah $F(t)$ bila

$$f(s) = \frac{s}{s^2 - k^2}$$

Jawab:

Fungsi $f(s)$ dapat diuraikan menjadi:

$$\begin{aligned} f(s) &= \frac{As + B}{s + k} + \frac{Cs + D}{s - k} \\ &= \frac{As^2 - Aks + Bs - Bk + Cs^2 + Ds + Cks + Dk}{s^2 - k^2} \end{aligned}$$

dengan substitusi balik, diperoleh $A = 0$, $B = 1/2$, $C = 0$, $D = 1/2$ sehingga



Turunan Transformasi Laplace

per-definisi:

$$L\{F'(t)\} = \int_0^{\infty} e^{-st} \frac{dF}{dt} dt$$

integrasi bagian:

$$\begin{aligned} L\{F'(t)\} &= e^{-st} F(t) \Big|_0^{\infty} + s \int_0^{\infty} e^{-st} F(t) dt \\ &= sL\{F(t)\} - F(0) \end{aligned}$$

kalau diteruskan

$$L\{F^{(2)}(t)\} = s^2 L\{F(t)\} - sF(t) - F'(0)$$

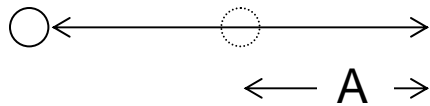
dan seterusnya:

$$L\{F^{(n)}(t)\} = s^n L\{F(t)\} - s^{n-1}F(t) - s^{n-1}F'(t) \dots - F^{(n-1)}(0)$$



Contoh di Fisika:

Kasus osilator harmonis:



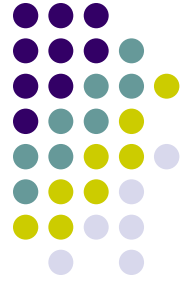
$$t=0, y = y_0, y' = 0$$

$$F = -ky \rightarrow y'' + \omega^2 y = 0$$

Kalau pada persamaan diferensial kita lakukan transformasi Laplace:

$$L\{y''\} = -\omega^2 L\{y\}$$

$$s^2 L\{y\} - sy(0) - y'(0) = -\omega^2 L\{y\}$$



masukkan syarat batas, diperoleh:

$$L\{y\} = \frac{s}{s^2 + \omega^2} y_0$$

inverse transformasi ini menghasilkan:

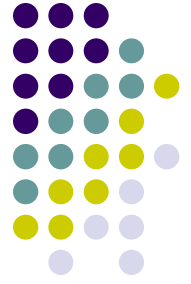
$$y = y_0 \cos \omega t \quad (\text{seperti yang diharapkan})$$

Pertanyaan, apabila syarat batas diubah menjadi

$$t=0, y=0, y'=v_0$$

apa yang terjadi?

(Jawab: $y = v_0/\omega \sin \omega t$, buktikan!)



Sifat-sifat lain fungsi Laplace

1. Substitusi

$$f(s-a) = L\{e^{at}F(t)\} \quad (\text{buktikan!})$$

Sehingga:

$$L\{e^{at} \sin kt\} = \frac{k}{(s-a)^2 + k^2}$$

$$L\{e^{at} \cos kt\} = \frac{(s-a)}{(s-a)^2 + k^2}$$

2. Translasi

$$e^{-bs} f(s) = L\{F(t-b)\} \quad (\text{buktikan!})$$

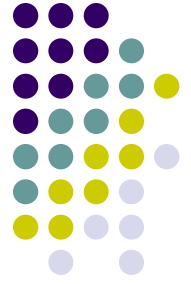
3. Turunan suatu transformasi

Turunan ke- n :

$$f^{(n)}(s) = L\{(-t)^n F(t)\} \quad (\text{buktikan!})$$

4. Integrasi suatu transformasi

$$\int_s^{\infty} f(x) dx = L\left\{\frac{F(t)}{t}\right\}$$



Contoh kasus: Osilator Teredam

Kasus getaran harmonis teredam:

$$mX''(t) + bX'(t) + kX(t) = 0$$

dengan m, k, b adalah konstan.

Bila kita gunakan kondisi inisial $X(0)=X_0$, $X'(0)=0$, maka persamaan transformasi menjadi:

$$m[s^2 x(s) - sX_0] + b[s x(s) - X_0] + kx(s) = 0$$

dan

$$x(s) = X_0 \frac{ms + b}{ms^2 + bs + k}$$



Persamaan terakhir dapat diselesaikan dengan melengkapi kuadrat penyebut sbb:

$$s^2 + \frac{b}{m}s + \frac{k}{m} = \left(s + \frac{b}{2m}\right)^2 + \left(\frac{k}{m} - \frac{b^2}{4m^2}\right)$$

Apabila faktor redaman (*damping*) kecil, $b^2 < 4km$, suku terakhir adalah positif dan sebut sebagai ω_1^2

$$\begin{aligned}x(s) &= X_0 \frac{s + b/m}{(s + b/2m)^2 + \omega_1^2} \\ &= X_0 \frac{s + b/2m}{(s + b/2m)^2 + \omega_1^2} + X_0 \frac{(b/2m\omega_1)\omega_1}{(s + b/2m)^2 + \omega_1^2}\end{aligned}$$



Kita gunakan

$$L\{e^{at} \sin kt\} = \frac{k}{(s-a)^2 + k^2}$$
$$L\{e^{at} \cos kt\} = \frac{(s-a)}{(s-a)^2 + k^2}$$

Didapat:

$$X(t) = X_0 e^{-(b/2m)t} \left(\cos \omega_1 t + \frac{b}{2m\omega_1} \sin \omega_1 t \right)$$

$$= X_0 \frac{\omega_0}{\omega_1} e^{-(b/2m)t} \cos(\omega_1 t - \varphi)$$

dengan

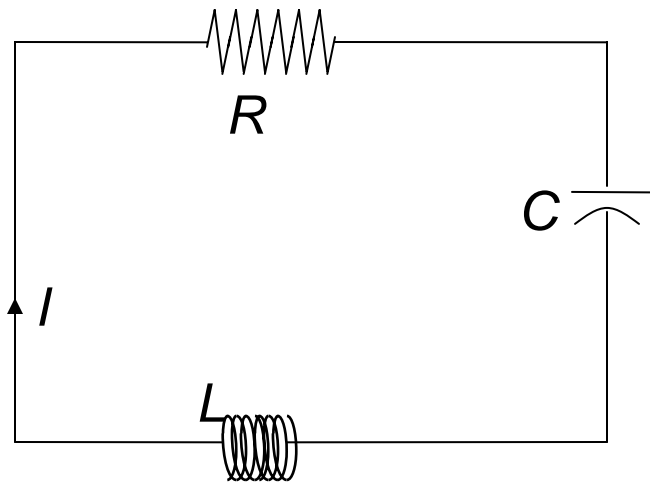
$$\tan \varphi = \frac{b}{2m\omega_1}$$

$$\omega_0^2 = \frac{k}{m}$$



RLC Analog

Ada keserupakan antara osilator harmonis teredam dengan rangkaian RLC



Dari hukum Kirchhoff:

$$L \frac{dI}{dt} + RI + \frac{q}{C} = 0$$

Didiferensialkan:

$$L \frac{d^2 I}{dt^2} + R \frac{dI}{dt} + \frac{1}{C} I = 0$$

$$mX''(t) + bX'(t) + kX(t) = 0$$

Analog dengan problem mekanika.



5.3. Transformasi Fourier

Secara Matematik transformasi Fourier dikembangkan dari deret Fourier. Secara detail dapat dilihat di Arfken.

$$g(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(t) e^{i\alpha t} dt$$

Inversnya:

$$f(t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(\alpha) e^{-i\alpha t} dt$$



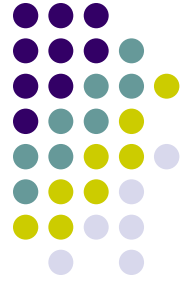
Berbagai macam bentuk TF

Pasangan transformasi Fourier

$$H(f) = \int_{-\infty}^{\infty} h(t)e^{-i2\pi ft} dt \Leftrightarrow h(t) = \int_{-\infty}^{\infty} H(f)e^{i2\pi ft} df$$

$$F(\alpha) = \int_{-\infty}^{\infty} f(x)e^{-i\alpha x} dx \Leftrightarrow f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} F(\alpha)e^{i\alpha x} d\alpha$$

$$g(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x)e^{ikx} dx \Leftrightarrow f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(k)e^{-ikx} dk$$



Di Mekanika Kuantum:

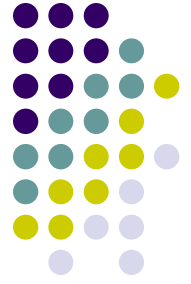
Paket gelombang, $f(x)$, dengan gelombang dalam bilangan gelombang, $g(k)$

$$f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(k) e^{-ikx} dk \quad \Leftrightarrow \quad g(k) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{ikx} dx$$

Lebih lengkap:

$$f(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} g(k) e^{-ik(x-\omega t)} dk$$

Jadi cukup banyak cara penulisan transformasi Fourier, pilih salah satu dan harus konsisten!



Sekarang kita lihat kenyataan bahwa pada umumnya hasil transformasi Fourier adalah fungsi kompleks.

$$H(f) = \int_{-\infty}^{\infty} h(t) e^{-2\pi i f t} dt$$

Kompleks

Maka:

$$\begin{aligned} H(f) &= R(f) + i I(f) \\ &\quad \text{(real)} \quad \text{(imaginer)} \\ &= |H(f)| e^{i\theta(f)} \end{aligned}$$

fase

dengan

$$|H(f)| = \sqrt{R^2(f) + I^2(f)}$$

$$\theta(f) = \arctan \left[\frac{I(f)}{R(f)} \right]$$



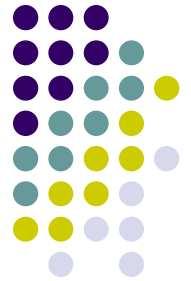
Contoh:

$$h(t) = \beta e^{-\alpha t} \quad \text{pada } t > 0 \\ = 0 \quad \text{pada } t < 0$$

Carilah harga real dan imajiner dari $H(f)$ dan fase $\theta(f)$

Jawab:

$$\begin{aligned} H(f) &= \int_0^{\infty} \beta e^{-\alpha t} e^{-i2\pi f t} dt \\ &= \int_0^{\infty} \beta e^{-(\alpha + i2\pi f)t} dt \\ &= \beta \cdot \frac{-1}{\alpha + i2\pi f} e^{-(\alpha + i2\pi f)t} \Big|_0^{\infty} \\ &= \frac{\beta}{\alpha + i2\pi f} \\ &= \frac{\alpha\beta}{\alpha^2 + (2\pi f)^2} - i \frac{\beta 2\pi f}{\alpha^2 + (2\pi f)^2} \end{aligned}$$

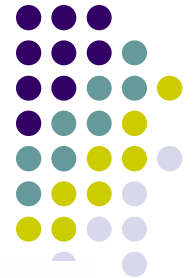


Sehingga:

$$R(f) = \frac{\alpha\beta}{\alpha^2 + (2\pi f)^2} \text{ dan } I(f) = \frac{\beta 2\pi f}{\alpha^2 + (2\pi f)^2}$$

serta

$$\theta(f) = \tan^{-1}\left(-\frac{2\pi f}{\alpha}\right) \text{ dan } |H(f)| = \frac{\beta}{\sqrt{\alpha^2 + (2\pi f)^2}}$$



Syarat keberadaan integral Fourier:

- 1). $\int_{-\infty}^{\infty} |h(t)| dt < \infty$ (finite) Syarat “cukup” tetapi tidak “harus”
→ maka $H(f)$ exists

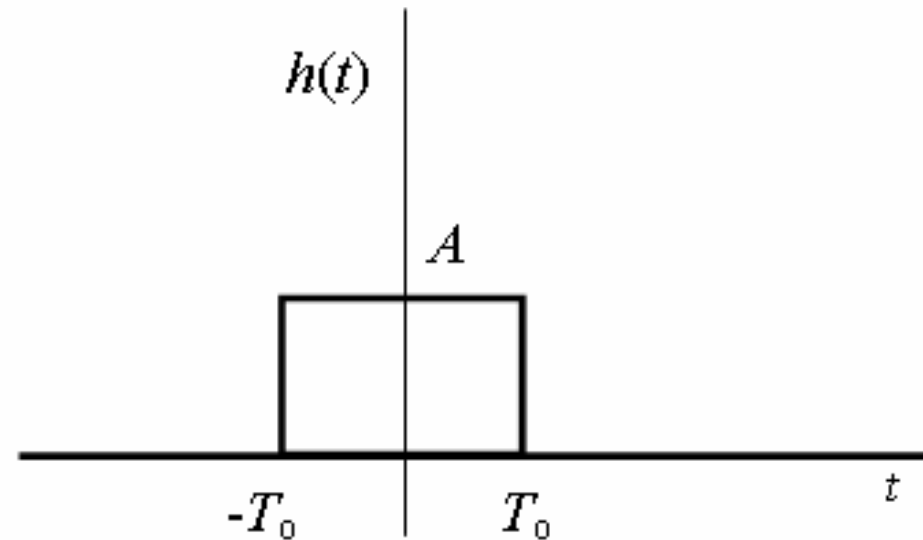
atau

- 2). Pada suatu fungsi
 $h(t) = \beta(t) \sin(2\pi\alpha t + \alpha)$ dengan α dan α konstanta sebarang
bila $\left| \frac{h(t)}{t} \right|$ integrable
→ maka $H(f)$ exists



Contoh-contoh

1. (untuk kondisi 1)
Gelombang pulsa

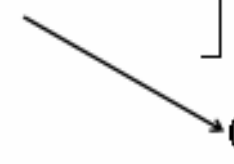


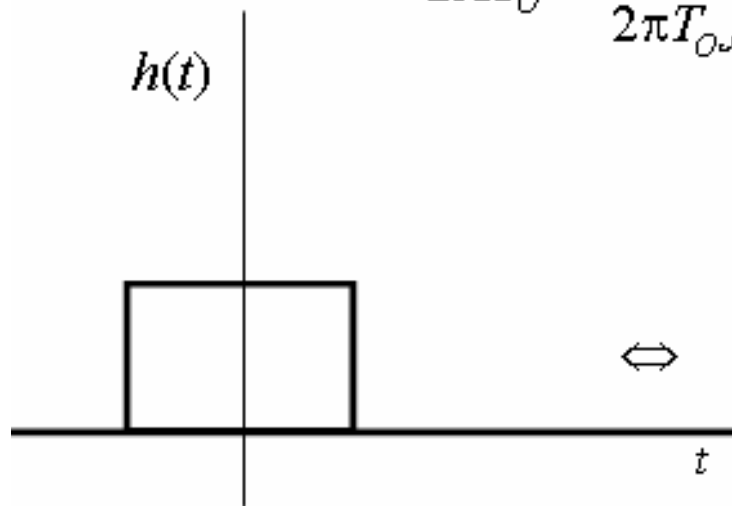
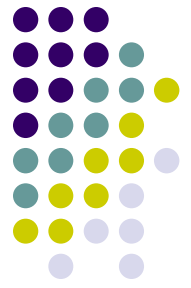
$$\begin{aligned} h(t) &= A & |t| < T_0 \\ &= 0 & |t| > T_0 \end{aligned}$$

$$\text{Kondisi 1: } \int_{-\infty}^{\infty} |h(t)| dt = \int_{-T_0}^{T_0} |A| dt = 2AT_0 \text{ (terpenuhi)}$$

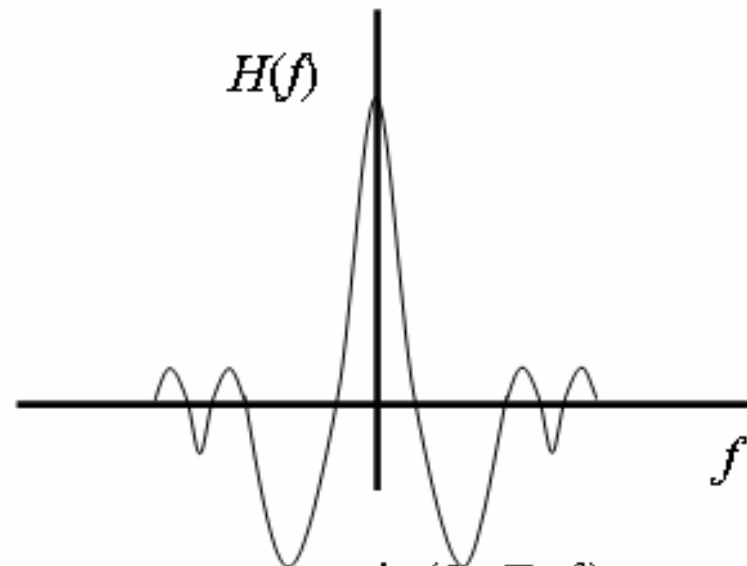
Jadi $H(f)$ exists, sekarang kita evaluasi:

$$\begin{aligned}
 H(f) &= \int_{-T_0}^{T_0} A e^{-i2\pi ft} dt \\
 &= A \left[\int_{-T_0}^{T_0} \cos 2\pi ft dt - i \int_{-T_0}^{T_0} \sin 2\pi ft dt \right] \\
 &= 2A \cdot \frac{1}{2\pi f} \sin 2\pi ft \Big|_0^{T_0} \\
 &= 2AT_0 \frac{\sin(2\pi T_0 f)}{2\pi T_0 f}
 \end{aligned}$$


 0 (ganjil)



$$h(t) = A \quad |t| < T_0 \quad \Leftrightarrow$$



$$H(f) = 2AT_0 \frac{\sin(2\pi T_0 f)}{2\pi T_0 f}$$



2. (untuk kondisi 2)

$$h(t) = 2Af_0 \frac{\sin(2\pi f_0 t)}{2\pi f_0 t}$$

dapat dibuktikan bahwa

$$\int_{-\infty}^{\infty} |h(t)| dt \text{ divergen (tak memenuhi kondisi 1)}$$

tetapi kondisi 2 masih terpenuhi, sehingga transformasi Fourier dapat dilakukan:

$$\begin{aligned} H(f) &= \int_{-\infty}^{\infty} h(t) e^{-2\pi i f t} dt \\ &= \frac{A}{\pi} \int_{-\infty}^{\infty} \frac{\sin(2\pi f_0 t)}{t} \{ \cos(2\pi f t) - i \sin(2\pi f t) \} dt \\ &= \frac{A}{\pi} \int_{-\infty}^{\infty} \frac{\sin(2\pi f_0 t) \cos(2\pi f t)}{t} dt \end{aligned}$$

ganjil

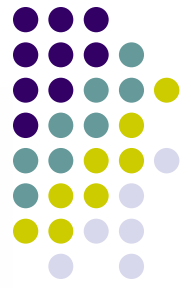


$$\begin{aligned} &= \frac{A}{2\pi} \int_{-\infty}^{\infty} \left\{ \frac{\sin 2\pi(f_0 + f)t}{t} + \frac{\sin 2\pi(f_0 - f)t}{t} \right\} dt \\ &= A \left\{ (f_0 + f) \int_{-\infty}^{\infty} \frac{\sin 2\pi(f_0 + f)t}{2\pi t} dt + (f_0 - f) \int_{-\infty}^{\infty} \frac{\sin 2\pi(f_0 - f)t}{2\pi t} dt \right\} \\ &= \frac{1}{2} A \left\{ (f_0 + f) \frac{1}{|f_0 + f|} + (f_0 - f) \frac{1}{|f_0 - f|} \right\} \end{aligned}$$

Pada $|f| < f_0$ disini $-f_0 < f < f_0$ $\rightarrow H(f) = A$
 $|f| = f_0$ $\rightarrow H(f) = A/2$
 $|f| > f_0$ disini $f > -f_0$ atau $f < f_0$ $\rightarrow H(f) = 0$

Dapat disimpulkan

$$h(t) = 2Af_0 \frac{\sin(2\pi f_0 t)}{2\pi f_0 t} \Leftrightarrow H(f) = A \quad |f| < f_0$$



3. Fungsi delta Dirac

$$h(t) = K\delta(t)$$

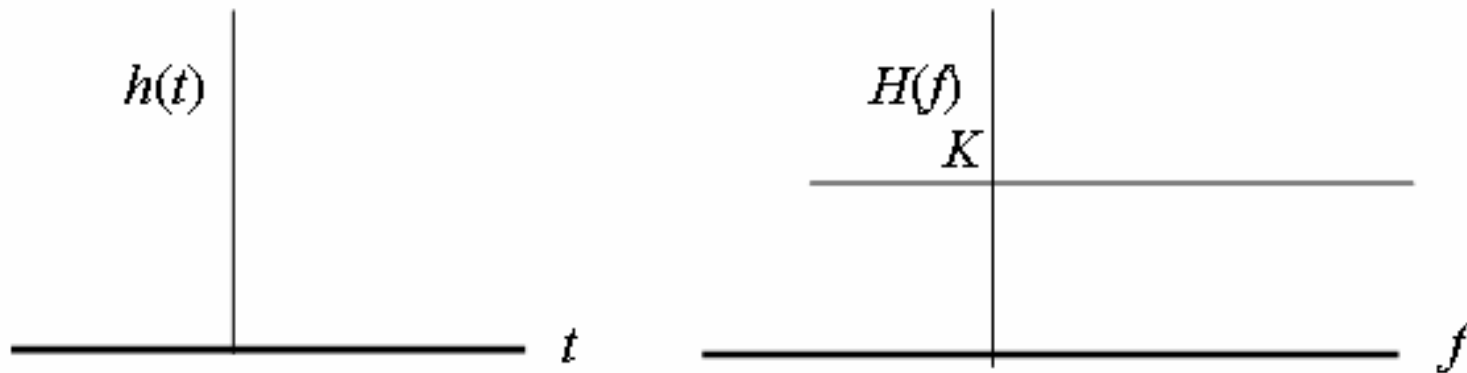
Hasil transformasi Fourier:

$$H(f) = \int_{-\infty}^{\infty} h(t)e^{-2\pi ift} dt = K \int_{-\infty}^{\infty} \delta(t)e^{-2\pi ift} dt = Ke^0 = K$$

Kebalikan:

$$h(t) = \int_{-\infty}^{\infty} H(f)e^{-2\pi ift} df = \int_{-\infty}^{\infty} Ke^{-2\pi ift} df = K\delta(t)$$

$$\text{Jadi: } h(t) = K\delta(t) \quad \Leftrightarrow \quad H(f) = K$$





4. Fungsi Periodik: (latihan)

$$h(t) = A \cos(2\pi f_o t)$$

Hasil transformasi Fourier:

$$\begin{aligned} H(f) &= \\ &= \\ &= \frac{A}{2} \delta(f - f_o) + \frac{A}{2} \delta(f + f_o) \end{aligned}$$

Topik-topik tersisa

- Teorema Konvolusi
- Representasi Momentum
- Persamaan Integral

